# Entropy of Faze Space of Physical Systems, Free and Bond Energy of Closed Physical Systems and their Relativity Properties 

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#### Abstract

In the article [1] we introduced the concept of entropy for such topological spaces that admit pseudo-convex coverings [1], and it was shown here that the class of such topological spaces is quite wide. The present article introduces the concepts of free and bond energies of a closed system, shows the relative nature this energies and entropy of the phase space of a closed system. There considered two case: the relative property are illustrated in two case a) when components of events are the coordinates of the vector, which length is equal to the total energy of the system in the isotropic basis. b) when components of events are the coordinates of the same vector in the orthonormal basis.


Keywords: Entropy; faze space; minkowski space; lorentz transformations.

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## 1 Introduction

Let's say there is a closed system with a phase space $X$ whose entropy is $S=\frac{N}{M}$ [1], where $N$ represents the number of elements in the minimal pseudo-convex open covering [1] of the topological space and $M$ the number of orbits acting on the covering elements of the homeomorphisms preserving this covering. For example, the entropy of a circle is three, and the entropy of a circle with a tail is $4 / 3$, The entropy of a sphere is 6 , and the entropy of a sphere with a tail is $6 / 3$. The concept of entropy for phase spaces of dynamical systems was used by us to describe the evolution of closed physical systems. By means of the concept of entropy, a random process was constructed that described the evolution of such systems in time [1,2,3]. Since the concept of entropy is the base in such constructions the behavior of entropy in different reference systems during these dynamic processes is interesting [4,5].

Free and bond energy of closed systems, Lorentz transformations of systems, Relative properties of free and bond energies Let's energy of closed system is $E$, call the quantity $E_{M}=\frac{\sqrt{E}}{S} M$ the bond energy of the system, and the quantity $E_{S}=\frac{\sqrt{E}}{M} S$ the free energy of the system. Obviously, that is $E=E_{M} E_{S}$, there is an inversely proportional relationship between the bond and free energies of the closed system.

In closed systems, entropy increases continuously at all moments of time, and the free energy of the system also increases. In such a system, time is projected into the change in the free energy of the system. Since every change in a closed system is the result of an increase in entropy, therefore, we can assume that the same that $E_{S}=C t$, where $t$ is time and $C$ is a speed of increase of free energy of systems which measured in $\mathrm{J} / \mathrm{sec}$. Lorentz transformation $L: R_{n-1}^{n} \rightarrow R_{n-1}^{n}$ of Minkowski $R_{n-1}^{n}$ space $[6,7,8,9]$ is transformation which saves scalar product: $\langle L(a), L(b)\rangle=\langle a, d\rangle$.

Consider the 2-dimensional vector space of events $R_{1}^{2}$, the elements of this space are pairs $\left(C t, E_{M}\right)$ of bond and free energies of closed systems. If we fix any isotropic base [10,9] in this space, and we will assume that the coordinates of event $\left(C t, E_{M}\right)[11,12,13]$ is given then in this basis, then square of thi pseudo-scalar product of each this event (vector) will be the product $E_{N} C t=E$.

For isotropic base Lorentz transformations [6,7,8,9] have the matrixes:

$$
\left(\begin{array}{ll}
a & 0 \\
0 & \frac{1}{a}
\end{array}\right)
$$

or

$$
\left(\begin{array}{ll}
0 & a \\
\frac{1}{a} & 0
\end{array}\right)
$$

Event $\left(C t, E_{M}\right)$ by this transformationsin event $\left(C t^{\prime}, E_{M}^{\prime}\right)$, where $E_{M}^{\prime}=a E_{M}+0 t, t^{\prime}=E_{M} 0+\frac{1}{a} t$. Indeed, $C t^{\prime}=\frac{1}{a} C t$ therefore $E^{\prime} a_{M} C t^{\prime}=E_{M} C \frac{1}{a} t=E_{M} C t=E$ for second matrix will be some.

$$
E^{\prime} a_{M} C t^{\prime}=E_{M} C \frac{1}{a} t=E_{M} C t=E
$$

Let's say now we have events in the system $\left(C t_{1}, E_{M}^{1}\right),\left(C t_{2}, E_{M}^{2}\right)$. This pair of events will have the form interval between these events [ $6,11,12,13$ ], it is

$$
s_{12}=\sqrt{\left(E_{M}^{2}-E_{M}^{1}\right)\left(C t_{2}-C t_{1}\right)}
$$

The interval between the corresponding two events $\left(C t_{1}^{\prime}, E_{M}^{\prime 1}\right),\left(C t_{2}^{\prime}, E_{M}^{2}\right)$ will be in the countdown system

$$
s_{12}^{\prime}=\sqrt{\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)\left(C t_{2}^{\prime}-C t_{1}^{\prime}\right)}
$$

If we make Lorentz transformation of one isotropic base to second, we will get

$$
E_{M}^{\prime 2}-E_{M}^{\prime 1}=a E_{M}^{2}-a E_{M}^{1}
$$

and

$$
C\left(t_{2}^{\prime}-t_{1}^{\prime}\right)=C\left(\frac{1}{a} t_{2}-\frac{1}{a} t_{1}\right)
$$

It follows from these last two equations

$$
s_{12}^{\prime}=\sqrt{a\left(E_{M}^{2}-E_{M}^{1}\right) \frac{C}{a}\left(t_{2}-t_{1}\right)}=s_{12}
$$

It means in two different isotropic reference system the energy and interval between event pairs are invariants.Therefore, the equality of the intervals between the events also took place in our case.

Consider the quantity $\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)$ using Lorentz transformations, we will have:

$$
E_{M}^{\prime 2}-E_{M}^{\prime 1}=a E_{M}^{2}-a E_{M}^{1}=a\left(E_{M}^{2}-E_{M}^{1}\right)
$$

It follows, that the measure of difference between bond energies in two events first isotropic referencesystem well be more than measure of difference between bond energies between two the respective events of second isotropic reference system if $a<1$. If $a>1$, it will be vice versa.

Consider new equation

$$
C\left(t_{2}^{\prime}-t_{1}^{\prime}\right)=C\left(\frac{1}{a} t_{2}-\frac{1}{a} t_{1}\right),
$$

From this equation follows, that, the measure of difference between free energies in two events for observer of first isotropic reference system will be less than measure of difference between free energies between two the respective events of second isotropic reference system if $a<1$. If $a>1$, it also will be vice versa.

Let new, consider orthonormal $e_{1}, e_{2}$ base inMinkowski vector space $R_{1}^{2}$. Let $K$ The reference system associated with this base, in such system for event $\left(C t, E_{M}\right)$ product $E_{M} C t$ It does not represent a scalar square of this event. In case of a orthonormal base we can consider second $K^{\prime}$ reference system obtained from the system $K$ by moving along the axis corresponding to component $E_{M}$ at $V$ joule/s a constant speed. In the isotropic case, we cannot consider such a reference system. Such transformations of coordinates in ortonormal case is called Lorentz t transformations [6,7,8,9,13] and has a form:

$$
\begin{aligned}
& E_{M}^{\prime}=\frac{E_{M}-V t}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}} \\
& t^{\prime}=\frac{t-\frac{V}{C^{2}} E_{M}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}
\end{aligned}
$$

As in the case of the isotropic basis, the energy of the system in the orthonormal basis is the square of the length of the vector representing the given event.

It is clear, that energy of system is invariant for Lorentz transformations, indeed

$$
\begin{aligned}
& E_{M}^{\prime 2}-C^{2} t^{\prime 2}=\frac{\left(E_{M}-V t\right)^{2}-C^{2}\left(\frac{V}{C^{2}} E_{M}-t\right)^{2}}{1-\left(\frac{V}{C}\right)^{2}}=\frac{E_{M}^{2}-2 E_{M} V t+V^{2} t^{2}-C^{2} \frac{V^{2}}{C^{4}} E_{M}^{2}+2 C^{2} E_{M} \frac{V}{C^{2}} t-C^{2} t^{2}}{1-\left(\frac{V}{C}\right)^{2}}= \\
& \frac{E_{M}^{2}\left(1-\left(\frac{V^{2}}{C^{2}}\right)\right)-C^{2} t^{2}\left(1-\frac{V^{2}}{C^{2}}\right)}{1-\left(\frac{V}{C}\right)^{2}}=E_{M}^{2}-C^{2} t^{2}
\end{aligned}
$$

In orthonormal reference system the interval between two pair of event $\left(C t_{1}, E_{M}^{1}\right),\left(C t_{2}, E_{M}^{2}\right) \quad$ in system $K$ is

$$
s_{12}=\sqrt{\left(C t_{2}-C t_{1}\right)^{2}-\left(E_{M}^{2}-E_{M}^{1}\right)^{2}}
$$

Interval is invariantfor such Lorentz transformation, Indeed, if $\left(C t_{1}^{\prime}, E_{M}^{\prime 1}\right),\left(C t_{2}^{\prime}, E_{M}^{\prime 2}\right)$ is pair of events insystem $K^{\prime}$ then interval between this events is

If we use the Lorentz transformation we well have:

$$
s_{12}^{\prime}=\sqrt{\left(C t_{2}^{\prime}-C^{\prime} t_{1}\right)^{2}-\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)^{2}}
$$

$$
C^{2}\left(t_{2}-t_{1}\right)^{2}=\frac{C^{2}\left(t_{2}^{\prime}-t_{1}^{\prime}\right)^{2}+2 V\left(t_{2}-t_{1}\right)\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)+\frac{V^{2}}{C^{2}}\left(E_{M}^{\prime 2}-E_{M}^{\prime \prime}\right)^{2}}{1-\frac{V^{2}}{C^{2}}}
$$

and

$$
\left(E_{M}^{2}-E_{M}^{1}\right)^{2}=\frac{\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)^{2}+2 V\left(t_{2}^{\prime}-t_{1}^{\prime}\right)\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)+V^{2}\left(t_{2}^{\prime}-t_{1}^{\prime}\right)^{2}}{1-\frac{V^{2}}{C^{2}}} .
$$

From this equations, finallywe will get:

$$
s_{12}=\sqrt{\left(C t_{2}-C t_{1}\right)^{2}-\left(E_{M}^{2}-E_{M}^{1}\right)^{2}}=s_{12}^{\prime}=\sqrt{\left(C t_{2}^{\prime}-C^{\prime} t_{1}\right)^{2}-\left(E_{M}^{2}-E_{M}^{\prime 1}\right)^{2}}
$$

Consider the magnitude $\sqrt{\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)^{2}}$, If we use the Lorentz transformation we well have:

$$
\sqrt{\left(E_{M}^{2}-E_{M}^{1}\right)^{2}}=E_{M}^{2}-E_{M}^{1}=\frac{E_{M}^{\prime 2}+V t^{\prime}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}-\frac{E_{M}^{\prime 1}+V t^{\prime}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}=\frac{E_{M}^{\prime 2}-E_{M}^{\prime 1}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}
$$

from this follows

$$
l=E_{M}^{\prime 2}-E_{M}^{\prime 1}=l_{0} \sqrt{1-\left(\frac{V}{C}\right)^{2}}
$$

where

$$
l_{0}=E_{M}^{2}-E_{M}^{1}
$$

This means that in stationary system, the difference in bond energies of two events in a moving system is smaller than for the stationary system.

The free energy in stationary reference system has form $E_{S}=C t$. If we take into account Lorentz transformations system, and the fact that the free energy in this moving system will have the form:

$$
C t^{\prime}=\frac{C t-C \frac{V}{C^{2}} E_{M}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}} .
$$

If we take into account, since in zero moment of time the reference systems $K, K^{\prime}$ are match and $E_{M}=0$ ,after time $t$ will be $E_{M}$ we can write $E_{M}=V t$ That's why we will have:

$$
C t^{\prime}=\frac{C t-C \frac{V}{C^{2}} E_{M}}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}=\frac{C t-\frac{V}{C^{2}} V C t}{\sqrt{1-\left(\frac{V}{C}\right)^{2}}}=C t \sqrt{1-\left(\frac{V}{C}\right)^{2}}
$$

For the magnitude we have $\sqrt{1-\left(\frac{V}{C}\right)^{2}} \leq 1$, therefore we can make conclusion: in stationary system the free energy grows more slowly than in moving system, hence the time also flows more slowly and entropy grows more slowly.

## 2 Conclusions

1. Based on the concept of entropy of the phase space of a closed physical system, we defined the concepts of free energy and bond energy of this system.
2. We represented free energy and bond energy as components of the vector-two-dimensional Minkowski space.
3. We showed: If components of events in $R_{1}^{2}$ Minkowski space is free energy and bond energy of closed physical space : $\left(E_{S}=C t, E_{M}\right)$ and $K, K^{\prime}$ two reference systems, for which $K$ stationary and $K^{\prime}$ moving whit speed $V$ joule/second, along axis appropriate to $E_{M}$, ten we have:
a) $s_{12}=\sqrt{\left(C t_{2}-C t_{1}\right)^{2}-\left(E_{M}^{2}-E_{M}^{1}\right)^{2}}=s_{12}^{\prime}=\sqrt{\left(C t_{2}^{\prime}-C^{\prime} t_{1}\right)^{2}-\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right)^{2}}$,
b) $E_{M}^{2}-E_{M}^{1}=\left(E_{M}^{\prime 2}-E_{M}^{\prime 1}\right) \sqrt{1-\left(\frac{V}{C}\right)^{2}}$,
c) $E_{s}^{\prime}=C t^{\prime}=\left(E_{s}=C t\right) \sqrt{1-\left(\frac{V}{C}\right)^{2}}$

The meaning of these equations is as follows:
Interval between pairs of free energy and bond of closed system is invariant for Lorentz transformation.
The different of projections of events $\left(C t_{1}, E_{M}^{1}\right),\left(C t_{2}, E_{M}^{2}\right)$ on the axis $E_{M}$ is more than The different of projections of events $\left(C^{\prime} t_{1}, E_{M}^{\prime 1}\right),\left(C t_{2}^{\prime}, E_{M}^{\prime 2}\right)$ on the axis $E_{M}^{\prime}$

The free energy of closed system in stationary reference system $K$ is more than in moving reference system $K^{\prime}$. This means, that free energy and entropy in system $K^{\prime}$, increase slowly than free energy and entropy in system $K$, it is the same as that time flows more slowly in a moving system $K^{\prime}$ than in a stationary system $K$ 4. The relevant dependencies are shown for that case when $K, K^{\prime}$ reference systems are isotropic and $K^{\prime}$ is obtained from $K$ by the above mentioned matrices.

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## Competing Interests

Author has declared that no competing interests exist.

## References

[1] Mumladze M. Entropy of Topological Space and Evolution of Phase Space of Dynamical Systems // International Journal of Management and Fuzzy Systems: 8-13.
Published Online: Jun. 15, 2020.
DOI: 10.11648/j.ijmfs. 2020060
[2] Fiodorov FI. Lorentz groups, Moscow. (in Russian). 1979;384c.
[3] Stochastic model of evolution of homotopy type of phase space of dynamical system,East European scientific Journal. 2021;14(68).
DOI: 10,32618/ESSA.2782-1994. 2021.1.68.17
[4] Entropy of State of Quantum System and Dynamic of This System, Open Access Library Journal. 2022;9(8).
[5] Ivo Terek Couto, Alexandre Lymberopoulos. Introduction to Lorentz Geometry, Curves and Surfaces, Imprint Chapman and Hall/CRC New York. 2021;350.
[6] Ivo Terek Couto, Alexandre Lymberopoulos. Introduction to Lorentz Geometry: Curves and Surfaces CRC press, Tailor \& Francis group, Carman \& Hall Book. 2021;340.
[7] Morínigo Fernando B, Wagner William, Pines David, Hatfield Brian. Feynman Lectures on Gravitation. West view Press. 2002;68. ISBN 978-0-8133-4038-8, Lecture 5.
[8] Pauli V. Теория Относительности, Moscow, (in Russian). 1947;300.
[9] Catoni F, et al. Mathematics of Minkowski Space. Frontiers in Mathematics. Basel: Birkhäuser Verlag; 2008.

DOI: 10.1007/978-3-7643-8614-6
ISBN 978-3-7643-8613- ISSN 1660-8046.
[10] Mumladze M. Basics of Geometry, Gori, (in Georgian). 2023;129.
[11] Dennery, Philippe; Krzywicki, André. Mathematics for Physicists. Courier Corporation; 2012. ISBN 978-0-486-15712-2.
[12] Gourgoulhon Eric. Special Relativity in General Frames: From Particles to Astrophysics. Springer. 2013;213. ISBN 978-3-642-37276-6.
[13] Koks, Don. Explorations in Mathematical Physics: The Concepts Behind an Elegant Language (illustrated ed.). Springer Science \& Business Media. 2006;234.
ISBN 978-0-387-32793-8. Extract of page 234.
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